# Applied Strain Effect on Superconducting Properties for Detwinned (Y, Gd)BCO Coated Conductors

Takumi Suzuki, Satoshi Awaji, Hidetoshi Oguro, and Kazuo Watanabe

Abstract-We succeeded in the complete detwin of (Y, Gd)BCO coated conductors using a newly developed in-situ tensile apparatus at high temperature. The intensities of (020) and (200) peaks in a longitudinal and a transverse directions of the tape measured by XRD were similar to each other for usual (Y, Gd)BCO coated conductors before annealing. However, the (200) peak disappears for the longitudinal direction after annealing under the tensile stress. It means that the coated conductors are detwinned by annealing under the tensile stress. We found that  $T_{\rm c}$  linearly decreases with increasing strain for the B-domain (b-axis). This behavior is similar with a uniaxial-pressure effect in (Y, Gd)BCO single crystals. However, the strain dependence of  $T_{\rm c}$  for the A-domain (a-axis) behaves as a power-law function, whereas the uniaxial-pressure dependency of  $T_{\rm c}$  for the a-axis is reported to be linear. In addition, the strain dependencies of  $J_{\rm c}$  are similar with those of  $T_c$ .

*Index Terms*—Coated conductor, detwin, strain effect, (Y, Gd)BCO.

## I. INTRODUCTION

HE REBa<sub>2</sub>Cu<sub>3</sub>O<sub>v</sub> coated conductors (REBCO, RE: rareearth element) have a large critical current density and strong mechanical properties. This high performance of the REBCO coated conductors enables us to develop the high field magnets operating at low temperatures. During an operation of a superconducting magnet, the coated conductor experiences various strains. Therefore the strain dependence of superconducting properties is a very important factor for superconducting magnet applications. The strain dependence of superconducting properties for REBCO coated conductors have been reported by many groups [1]–[5]. However, the strain effects of REBCO coated conductors are too complicated to be understood because of the existence of twin boundaries. This means that the a- and b-axis domains are mixed and aligned along the longitudinal direction of the tape. The volume fraction of the a- and b-axis domains in the longitudinal direction of usual conductors is about 50% [5].

On the other hand, it is reported that the  $T_c$  of a single crystal has a linear stress dependence [6]. The uniaxial pressure dependence of  $T_c$  for the *a*- and *b*- axes are opposite each other [6], [7]. However, the uniaxial strain dependence of  $J_c$ 

The authors are with the Institute for Materials Research, Tohoku University, Sendai 980-8577, Japan (e-mail: takumi-s@imr.tohoku.ac.jp).

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for the biaxially textured coated conductors in a self-field obeys a power-law function [2], [3]. Therefore, it is difficult to infer the strain dependence of  $J_c$  from that of  $T_c$ . In particular, the double peaks appear in the strain dependence of  $J_c$  in a low field region and disappear in a high field region above 3T [4]. The mixing of the *a*- and *b*-axis domains may be one of the origins for the complex strain dependence of  $J_c$  in the REBCO coated conductors.

Recently, the strain dependence of  $J_c$  is described by the 2D internal strain model, assuming the strain sensitivity for each crystal axis [5]. The internal strain state of a GdBCO layer in a coated conductor was evaluated at RT by a diffraction technique using synchrotron radiation. From the obtained results of internal strains, the strain sensitivities of  $J_c$  for the *a*- and *b*- axes were estimated using the strain sensitivities, assuming that they both obey a power-law function. Positive and negative strain sensitivities were derived for the *a*- and *b*-axes from this calculation. From these analyses, it is required strongly to determine the strain dependence of  $T_c$  and  $J_c$  the each axis for a wide strain range experimentally.

In case of (Y, Gd)BCO single crystal, the twin boundary can be disappeared by annealing under the strain [8]. We have tried to control the twin boundaries and residual strain by the strainannealing for the REBCO coated conductors using a bending jig [9], [10]. However, this method is not enough to remove twin boundaries for coated conductors completely, because of a formation of cracks. In this study, we succeeded in preparing the completely detwined REBCO coated conductors by a newly developed in-situ tensile apparatus at high temperature. We also investigate the single-crystal-like uniaxial strain dependence of superconducting properties of REBCO coated conductors. The applied strain effects on  $T_c$  and  $J_c$  for detwinned (Y, Gd)BCO coated conductors at 77 K are reported in this paper.

## **II. EXPERIMENTAL DETAIL**

Samples used in this study were a commercially available (Y, Gd)BCO coated conductor without an artificial pinning center made by the SuperPower Inc [11]. The substrate is a biaxially textured buffer made by the ion beam assisted deposition (IBAD) of MgO on a high-strength Hastelloy tape with a thickness of 100  $\mu$ m. The (Y, Gd)BCO layer was grown by Metal Organic Chemical Vapor Deposition (MOCVD). The annealing under the strain is carried out by a newly developed in-situ tensile apparatus which tensile device is set in the image furnace. The samples were annealed at 300 °C for 3 hours in air with an external stress/strain. The tensile stress was gradually applied after the sample temperature reached to 300 °C and

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Fig. 1. Strain-stress curve for (Y, Gd)BCO coated conductors during the detwin process.

kept a constant displacement during annealing. However, the stress and strain increase when temperature decreases due to the thermal contraction. We controlled the applied strain less than about 0.9 % as shown Fig. 1, because an expansion of about 0.85 % is expected by the alignments of all the *b*-axis. Then the applied strain was released at 200 °C. The stress and strain were measured using a load cell and a strain gauge, respectively.

The lattice constants and crystal alignments of the a- and b-axes along the longitudinal and transverse directions of the (Y, Gd)BCO coated conductors were measured at room temperature (RT) by a transmission X-ray diffraction using MoK $\alpha$  X-ray. The strain dependence of  $T_c$  and  $J_c$  were measured using a CuBe beam placed in a four-point bending apparatus [12]. The  $J_c$  measurements under the strain were performed at 77 K in liquid N<sub>2</sub>. The samples were soldered onto the surface of the beam. The transport properties were measured by a four probe method under the applied strain from a compressive to a tensile strain. We define a compressive strain as a negative value in this study. A micro-bridge with about 100, 200 or 500  $\mu m$  wide and 1 mm long was made by a photolithography. The  $J_{\rm c}$  was determined using the criterion of  $E = 1 \ \mu \text{V/cm}$ . Strains were measured using a strain gauge on the samples.

# III. EXPERIMENTAL RESULT

Fig. 2 shows the X-ray diffraction patterns along the longitudinal and transverse directions of the (Y, Gd)BCO coated conductors. Both (020) and (200) reflections can be seen. The amount of the *a*- and *b*-axis domains in the longitudinal direction of the as-received (Y, Gd)BCO tape is about 50%, because the intensities of (020) and (200) peaks are similar to each other. However, the annealing under the strain changes the peak intensity and position at the same time. The (200) peak in the longitudinal direction disappears after annealing under the strain. It means that the coated conductors can be detwinned by annealing under the strain. From the changes of the intensity ratio, we found that the domain ratios for the longitudinal and transverse directions become  $V_{\rm a}: V_{\rm b} = 94$ : 6 and = 6:94, respectively. Here we call the longitudinal and transverse directions of the strain-annealed samples A-domain and B-domain samples, respectively. Therefore, the A-domain



Fig. 2. (200) and (020) reflections along the longitudinal and transverse directions after the annealing under the strains in the (Y, Gd)BCO coated conductors. The longitudinal direction data of the as-received sample are also shown for comparison. Intensities are normalized by the maximum values.

and B-domain samples were taken out from the same sample with different direction.

From the changes of peak position, we calculated the effective residual strain along the longitudinal direction of the tapes. The relative change of residual strain was estimated using the difference of *d*-values from the as-received sample. Hence, the strain of the as-received sample was assumed to be zero. The relative residual strain can be defined as follows:

$$\varepsilon_{(hkl)}^{*} = \frac{d_{(hkl)} - d_{0(hkl)}}{d_{0(hkl)}}$$
(1)

where  $\varepsilon^*_{(hkl)}$ ,  $d_{(hkl)}$ , and  $d_{0(hkl)}$  are the effective residual lattice strain, the lattice spacing of the sample and the as-received sample in terms of the (hkl) reflection, respectively. The  $\varepsilon^*_{(200)}$  of A-domain and  $\varepsilon^*_{(020)}$  of B-domain samples becomes -0.066%and 0.227%, respectively.

Fig. 3 shows a typical example of the SEM surface images after the strain annealing with a maximum strain of about 0.9%. Many micro-cracks exist in the entirety of sample. The size of micro cracks is approximately  $0.1 \times 200 \ \mu\text{m}^2$ . Most of cracks align perpendicular to the applied strain direction, i.e., along the *a*-axis. Therefore the  $J_c$  was deteriorated due to the crack. The  $J_c$  values at 77 K without applied strain are 2.6, 0.79, and 0.013 MA/cm<sup>2</sup> for the as-received, the A-domain and the B-domain samples respectively.

The change in lattice strain was evaluated under uniaxial tensile loading of the B-domain sample. Fig. 4 shows the applied strain dependence of lattice strain along longitudinal direction of the B-domain sample. The lattice strain was determined from equation (1). In this case, the  $d_0$  is the lattice spacing of the B-domain sample at 0%. The lattice strain increases along the *b*-axis of the B-domain sample with increasing applied strain. The lattice strain of the B-domain sample has linear applied strain dependence. The slope of the applied strain dependence of lattice strain (strain ratio) for (Y, Gd)BCO (020) and Hastelloy become 0.57 and 0.88 respectively. On the other hand, the difference between the lattice strain and the applied strain along the longitudinal direction of the REBCO coated conductor was reported [5], [13]. Although the strain ratio obtained in this study is smaller than the strain ratio reported for



Fig. 3. SEM surface image of (Y, Gd)BCO layer after the annealing under the strains. (a) low magnification, (b) high magnification.



Fig. 4. Applied strain dependence of lattice strain after annealing under the strains in the (Y, Gd)BCO coated conductors along longitudinal direction.

the similar REBCO coated conductors made by the SuperPower in ref. [14], the linear applied strain dependence of lattice strain is observed. Therefore, it can be considered that the effect of the micro-cracks on the strain dependence is related to only the strain ratio.

Fig. 5 shows the strain dependences of  $T_c$ . The  $T_c$  for the B-domain (*b*-axis) decreases, but for the A-domain (*a*-axis) increases when the applied strain increases. The strain dependence of  $T_c$  for the A-domain (*a*-axis) can be well fitted by a power-law function whereas the single crystal behavior is linear [6], [7]. It was impossible to measure the pressure effect with a wide pressure region, since a brittle single crystal is easily fractured by the high pressure. Actually, the maximum pressure reported in ref. [7] corresponds to about 0.07%,



Fig. 5. Strain dependence of  $T_c(\varepsilon)/T_c(\varepsilon = 0\%)$  after annealing under the strains in the (Y, Gd)BCO coated conductors. The A-domain and B-domain mean the applied strain along a-axis and b-axis directions, respectively.  $T_c(+0\%)$  and  $T_c(-0\%)$  is  $T_c$  (0%) for measurement of tensile region and compressive region respectively.

which is calculated using a Young's modulus of 162.7 GPa for a-axis [14]. Therefore we found that the strain dependence of  $T_c$  for A-domain shows the power-law behavior from the strain dependence measurement for a wide strain region, which was about 10 times wider than that reported in the single crystal. On the other hand, the strain dependence of  $T_c$  for the A-domain was analyzed by using the following empirical function:

$$\frac{T_c(\varepsilon)}{T_c(0\%)} = \frac{T_{c,\max}}{T_c(0\%;)} - a'(\varepsilon - \varepsilon_p)^2$$
(2)

where a' is a strain sensitivity for the  $T_c$ . The maximum strain  $\varepsilon_p$ , where the  $T_c$  becomes a maximum, is estimated to be 0.77%. In addition, the additional thermal strain caused from the difference of thermal expansion between (Y, Gd)BCO and Cu-Be induces the compression strain of about 0.15% to the (Y, Gd)BCO. Hence, the strain at the maximum of  $T_c$  is shifted to a tensile strain side by the thermal strain. However, the  $\varepsilon_p$  value is much higher than the residual strain of 0.15%. Hence, the physical meaning of the  $\varepsilon_p$  is not clear yet. It should be discussed on the basis of the mechanism of high- $T_c$  superconductors.

Fig. 6 shows the strain dependencies of normalized  $J_c$  at 77.3 K. The strain dependencies of  $J_c$  for the as-received and the A-domain samples behave as a power-law function similarly to that of  $T_c$  and the B-domain sample is linear. The strain dependency was analyzed by using the following empirical function [5]:

$$\frac{J_c(\varepsilon)}{J_c(0\%)} = \frac{J_{c,\max}}{J_c(0\%)} - a(\varepsilon - \varepsilon_p)^2$$
(3a)

$$\frac{J_c(\varepsilon)}{J_c(0\%)} = b - a\varepsilon \tag{3b}$$

where a is a strain sensitivity for the  $J_c$ , b is a fitting parameter and  $\varepsilon_p$  is a maximum strain, at which the  $I_c/I_c$  (0%) becomes the maximum. The  $\varepsilon_p$ -values for the A-domain and the asreceived samples are 0.249% and 0.006% respectively. It is noticed that the  $\varepsilon_p$  values are much different between  $T_c$  and  $J_c$  for the a-axis. The strain sensitivity a-values for the A-domain and



Fig. 6. Strain dependence of  $J_c/J_c(\varepsilon = 0\%)$  after the annealing under the strains and as-received in the (Y, Gd)BCO coated conductors. The A-domain and B-domain are longitudinal and transverse direction respectively.



Fig. 7. Lattice strain dependence of  $J_c/J_c(\varepsilon = 0\%)$  after the annealing under the strains and as-received in the (Y, Gd)BCO coated conductors. A-domain and B-domain are axial and lateral direction, respectively.

the as-received samples are 0.30 and 0.17 respectively. In case of the B-domain sample, however, the  $J_c$  linearly decreases with increasing applied strain. Therefore, it is considered the strain dependencies of  $J_c$  are related to those of  $T_c$ .

Fig. 7 shows the lattice strain dependence of  $J_c(\varepsilon)/J_c(0\%)$ . The strain ratio of the *b*-axis for B-domain sample is 0.57. For the A-domain sample, because the applied strain direction is parallel to the micro-crack direction, it can be considered that the effect of micro-crack on the strain ratio is neglected for the a-axis. Hence, we use the strain ratio of 0.88 along the *a*-axis for the A-domain sample, which was reported in ref. 13. The lattice strain sensitivities for the *a*- and *b*-axes are 0.38 and 0.19, respectively from Fig. 7. It is concluded that the lattice strain sensitivity for each axis obtained in this study is an intrinsic strain sensitivity.

## **IV. CONCLUSION**

We succeeded in the preparation of completely detwinned (Y, Gd)BCO coated conductors using a newly developed high temperature tensile apparatus. Applied strain effects on  $T_c$  and  $J_c$  for the detwinned (Y, Gd)BCO coated conductors were also investigated. The intensities of (020) and (200) peaks along the longitudinal directions measured by XRD were similar to

each other for (Y, Gd)BCO coated conductors before annealing. However, the (200) peak in the longitudinal direction disappears after annealing under the tensile stress. It means that the coated conductors are detwinned by annealing under the tensile stress. From the XRD for (020) and (200) peaks of (Y, Gd)BCO, we confirmed that the domain ratio of the aand b-axes aligned along the longitudinal direction of the (Y, Gd)BCO tape becomes 94:6. Although the micro-cracks perpendicular to the longitudinal direction exist in the entirety of sample, we confirmed that internal strain increase with increasing applied strain by XRD. We found that  $T_{\rm c}$  linearly decreases with increasing strain for the B-domain but obeys a power-law for the A-domain. The strain dependences of  $J_{\rm c}$  are similar with those of  $T_{\rm c}$ . This means that the strain dependence of  $J_{\rm c}$  is closely related to that of  $T_{\rm c}$ . The lattice strain sensitivities of  $J_c$  for a- and b-axes are estimated to be 0.38 and 0.19, respectively.

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